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# An anomalous dip in thermoelectric power of $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$

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## Abstract

The thermoelectric power,  $S$ , has been studied for sintered samples of  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$  with  $0 \leq x \leq 0.30$  in the temperature range from the superconducting  $T_c$  to room temperature.  $S$  increases with decreasing temperature, and has a broad maximum at  $T^{\text{max}}$  in the region around 120 K before decreasing strongly when  $T_c$  is approached. Several properties indicate a decrease of charge concentration with increasing doping,  $x$ , from  $S(x, 290 \text{ K})$ ,  $T^{\text{max}}$ , and the resistivity  $\rho(x, 290 \text{ K})$ . An anomaly has been observed in  $S(T)$  for  $x \geq 0.20$  in the form of a dip at 78 K of order 15% of  $S$ . The origin of this feature is not known.

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*Keywords:* Superconductivity; Thermoelectric power; Pr doping

## 1. Introduction

In the RE-123 (RE = rare earth element) superconductors the RE's can be exchanged for each other with only small effect on the critical temperature  $T_c$  with the exception of Pr substitution which decreases  $T_c$  dramatically. Different explanations [1–4] for the depression of  $T_c$  have been suggested. Varying results for the valence of Pr is one reason. Recent neutron diffraction results in Nd(Pr)-123 [5] and analysis of  $T_c$  in Y(Pr)-123 [6] have suggested that hole localisation in the  $\text{Pr}^{4+}$  site is the main reason for the suppression of superconductivity by Pr. The observation of superconductivity in single crystals of  $\text{PrBa}_2\text{Cu}_3\text{O}_{7-\delta}$  [7,8] further complicates the understanding of the role played by Pr in the suppression of  $T_c$ . In our previous work [5], resistivity measurements in  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_7$  showed a decreasing metallic behaviour with increased resistivity and decreased slope of the normal state resistivity vs. temperature with increasing  $x$ . Neutron diffraction data

indicated that the Cu1–O4 and Cu2–O4 distances and the oxygen content were independent of Pr doping in Nd(Pr)-123. Bond valence sum (BVS) calculations showed a constant Cu2 valence but a decrease of the total hole concentration in the  $\text{CuO}_2$  plane. To further investigate the effect of Pr in  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$  system, we have measured the thermoelectric power  $S(x, T)$ .

## 2. Sample characterisation and experimental

Samples of  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$  (with  $x = 0, 0.05, 0.10, 0.15, 0.20$  and  $0.30$ ) were prepared by standard solid-state methods. Starting materials were high purity  $\text{Nd}_2\text{O}_3$ ,  $\text{BaCO}_3$ ,  $\text{CuO}$ , and  $\text{Pr}_6\text{O}_{11}$ . The samples were pressed into pellets and calcinated in air at 900, 920 and 920 °C with intermediate grindings. They were then annealed in flowing oxygen at 460 °C for three days and the temperature was finally decreased to room temperature at a rate of 12 °C/h.

The samples were characterised by X-ray powder diffraction (XRD). The XRD patterns were recorded in a Guinier–Hägg focusing camera using  $\text{CuK}_\alpha$  radiation with Si as an internal standard [9]. The XRD results for

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Pr-doped Nd-123 samples displayed single-phase behaviour. All XRD patterns were indexed with an orthorhombic unit cell.

The electrical resistivity was measured with a standard dc four-probe method. Electrical leads were attached to the sample by silver paint and heat treated at 300 °C in flowing oxygen for half an hour, which gave contact resistances of order 1–2  $\Omega$ . Thermoelectric power measurements were made on sintered bars of typical dimensions  $0.5 \times 2.5 \times 10 \text{ mm}^3$ , using a small, reversible temperature difference of 1.5 K.

### 3. Results and discussion

Fig. 1 shows the thermoelectric power  $S$  as a function of temperature and Pr concentration for  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$ .  $S$  first increases with increasing temperature towards a broad maximum at  $T^{\text{max}}$  above the superconducting  $T_c$ , and then decreases up to room temperature. Both  $S$  and  $T^{\text{max}}$  increase with increasing  $x$ . As further illustrated in the inset of Fig. 1, a dip in  $S(T)$  was observed at a temperature in the range 78–79 K. This feature only occurred for  $x \geq 0.20$ . The proximity to the boiling point of liquid nitrogen might suggest an experimental artefact from nitrogen boil off. Such an effect appears to be unlikely however, since reproducible dips in  $S$  were observed in repeated experiments only for  $x \geq 0.20$  and not for the other samples. The origin of the anomaly is not known. Except for the dips, all studied properties of the samples with  $x \geq 0.20$  (e.g.  $S(290 \text{ K})$ ,  $\rho(290 \text{ K})$ ,  $T_c$ ,  $T^{\text{max}}$ , and BVS calculations) are in agreement with the trends expected from the results for  $x < 0.20$ .

The doping concentration dependence of the room temperature  $S(290 \text{ K})$  and resistivity  $\rho(290 \text{ K})$  are

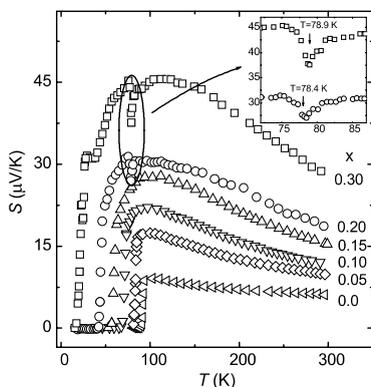


Fig. 1. Thermoelectric power  $S$  as a function of temperature and Pr doping for  $\text{Nd}_{1-x}\text{Pr}_x\text{Ba}_2\text{Cu}_3\text{O}_{7-\delta}$ . Inset:  $S(T)$  in the temperature range 72–87 K, where an anomaly has been observed in the form a dip, for  $x = 0.20$  and  $0.30$ .

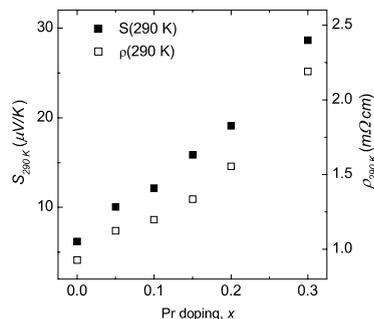


Fig. 2.  $S(290 \text{ K})$  (left scale) and  $\rho(290 \text{ K})$  (right scale) vs. Pr doping,  $x$ .

plotted in Fig. 2. Both  $S(290 \text{ K})$  and  $\rho(290 \text{ K})$  show a continuous increase with increasing Pr doping. The results for  $S(290 \text{ K})$ ,  $T^{\text{max}}$ , and  $\rho(290 \text{ K})$  suggest that Pr reduces hole concentration in the  $\text{CuO}_2$  plane. This is well in agreement with BVS calculations from neutron diffraction data [5].

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